

Al-Ayen University
College of Petroleum
Engineering

Numerical Methods and Reservoir Simulation

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Lecture 5: Principles of Finite Difference Approximation (Finite Difference Scheme & Discretization)

Finite Differences

Definition: The finite difference scheme is a way of approximating derivatives of a function. There are three ways to do that :

Approximation 1 -Forward Differences

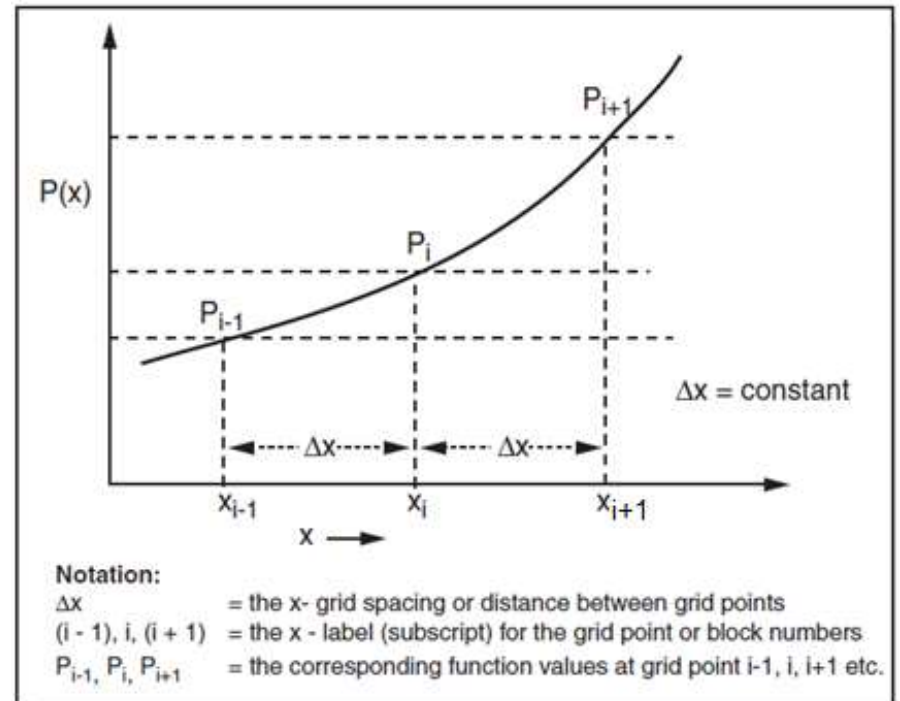
(fd):
$$\left(\frac{dP}{dx}\right)_{ifd} = \frac{P_{i+1} - P_i}{\Delta x}$$

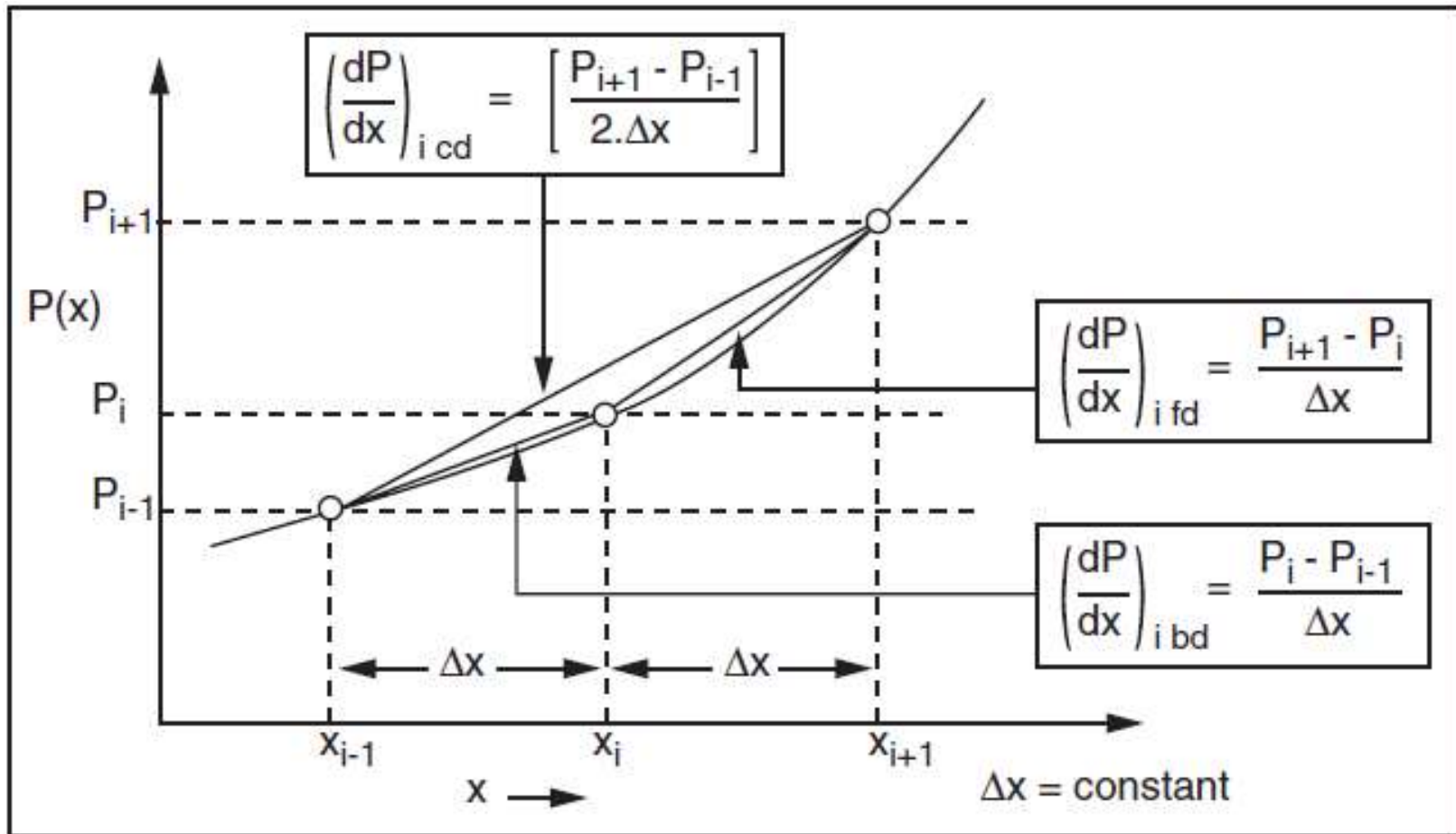
Approximation 2 -Backward Differences

(bd):
$$\left(\frac{dP}{dx}\right)_{ibd} = \frac{P_i - P_{i-1}}{\Delta x}$$

Approximation 3 -Central Differences

(cd):
$$\left(\frac{dP}{dx}\right)_{icd} = \frac{1}{2} \left[\left(\frac{dP}{dx}\right)_{ifd} + \left(\frac{dP}{dx}\right)_{ibd} \right] = \frac{1}{2} \left[\frac{P_{i+1} - P_i}{\Delta x} + \frac{P_i - P_{i-1}}{\Delta x} \right]$$
$$\left(\frac{dP}{dx}\right)_{icd} = \left[\frac{P_{i+1} - P_{i-1}}{2 \cdot \Delta x} \right]$$





Graphical illustration of the finite difference derivatives calculated by backward differences (*bd*), forward differences (*fd*) and central differences(*cd*) .

- **Now consider the finite difference approximation of the second derivative:**

$$\left(\frac{\partial^2 P}{\partial x^2}\right) \approx \frac{\left(\frac{dP}{dx}\right)_{i\,fd} - \left(\frac{dP}{dx}\right)_{i\,bd}}{\Delta x}$$

$$\approx \frac{\left(\frac{P_{i+1} - P_i}{\Delta x}\right)_{i\,fd} - \left(\frac{P_i - P_{i-1}}{\Delta x}\right)_{i\,bd}}{\Delta x}$$

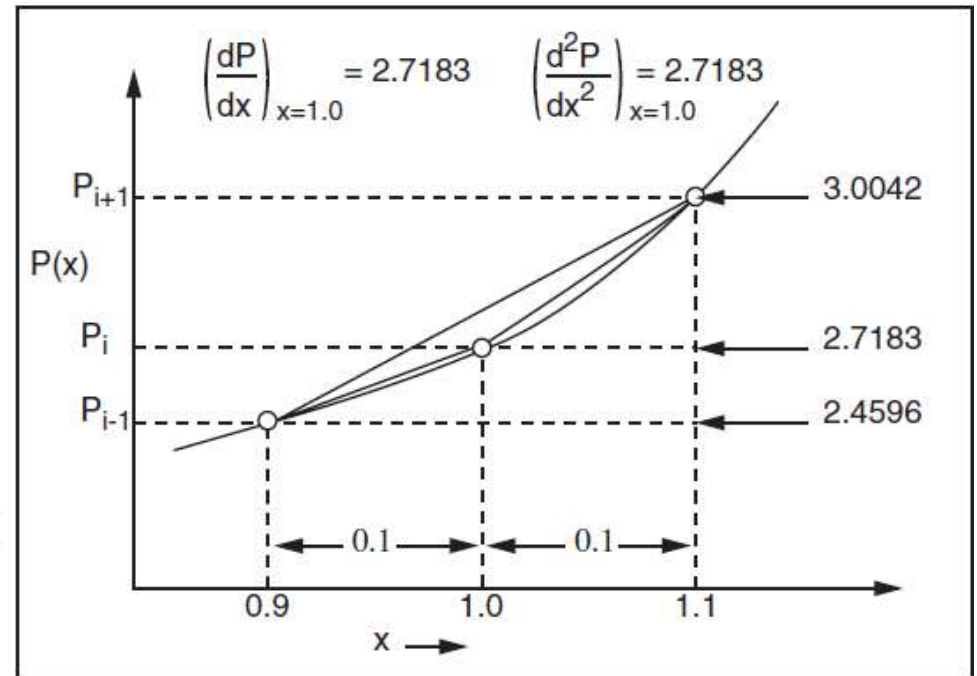
$$\left(\frac{\partial^2 P}{\partial x^2}\right) \approx \frac{(P_{i+1} + P_{i-1} - 2P_i)}{\Delta x^2}$$

Numerical Example

$$\left(\frac{dP}{dx}\right)_{ifd} \approx \left[\frac{3.0042 - 2.7183}{0.1}\right] = 2.859 \quad [err \approx +0.14]$$

$$\left(\frac{dP}{dx}\right)_{ibd} \approx \left[\frac{2.7183 - 2.4596}{0.1}\right] = 2.587 \quad [err \approx -0.13]$$

$$\left(\frac{dP}{dx}\right)_{icd} \approx \left[\frac{3.0042 - 2.4596}{2 \times 0.1}\right] = 2.723 \quad [err \approx -0.005]$$



A numerical example where the function values and derivative values are known ($P(x) = e^x$)

$$\left(\frac{\partial^2 P}{\partial x^2}\right) \approx \frac{3.0042 + 2.4596 - 2 \times 2.7183}{0.1^2} = 2.7200 \quad [err \approx 0.0017]$$

Identifying Error of Finite Difference Using Taylor Series

- A so-called Taylor series approximation of a function $f(x+h)$ expressed in terms of $f(x)$ and its derivatives $f'(x)$ may be written:

$$f(x+h) = f(x) + \frac{h}{1!} f'(x) + \frac{h^2}{2!} f''(x) + \frac{h^3}{3!} f'''(x) + \dots$$

- Applying Taylor series to our pressure function.

Approximation of the second order space derivative

At constant time, t , the pressure function may be expanded forward and backwards:

$$P(x + \Delta x, t) = P(x, t) + \frac{\Delta x}{1!} P'(x, t) + \frac{(\Delta x)^2}{2!} P''(x, t) + \frac{(\Delta x)^3}{3!} P'''(x, t) + \dots$$

$$P(x - \Delta x, t) = P(x, t) + \frac{(-\Delta x)}{1!} P'(x, t) + \frac{(-\Delta x)^2}{2!} P''(x, t) + \frac{(-\Delta x)^3}{3!} P'''(x, t) + \dots$$

By adding these two expressions, and solving for the second derivative, we get the following approximation:

$$P''(x, t) = \frac{P(x + \Delta x, t) - 2P(x, t) + P(x - \Delta x, t)}{(\Delta x)^2} + \frac{(\Delta x)^2}{12} P''''(x, t) + \dots$$

$$P''(x,t) = \frac{P(x + \Delta x, t) - 2P(x, t) + P(x - \Delta x, t)}{(\Delta x)^2} + \frac{(\Delta x)^2}{12} P''''(x, t) + \dots$$

or, by employing the grid index system, and using superscript to indicate time level:

$$\left(\frac{\partial^2 P}{\partial x^2}\right)_i^t = \frac{P_{i+1}^t - 2P_i^t + P_{i-1}^t}{(\Delta x)^2} + O(\Delta x^2)$$

- Here, the rest of the terms from the Taylor series expansion are collectively denoted $O(\Delta x^2)$ thus denoting that they are in order of, or proportional in size to Δx^2
- This error term is called **discretization error (or truncation error)**, which in this case is of second order, is neglected in the numerical solution.
- The smaller the grid blocks used, the smaller will be the error involved.
- Any time level could be used in the expansions above.

$$\left(\frac{\partial^2 P}{\partial x^2}\right)_i^{t+\Delta t} = \frac{P_{i+1}^{t+\Delta t} - 2P_i^{t+\Delta t} + P_{i-1}^{t+\Delta t}}{(\Delta x)^2} + O(\Delta x^2)$$

$$\left(\frac{\partial^2 P}{\partial x^2}\right)_i^{t+\frac{\Delta t}{2}} = \frac{P_{i+1}^{t+\frac{\Delta t}{2}} - 2P_i^{t+\frac{\Delta t}{2}} + P_{i-1}^{t+\frac{\Delta t}{2}}}{(\Delta x)^2} + O(\Delta x^2)$$

Approximation of the time derivative

At constant position, x , the pressure function may be expanded in forward direction in regard to time:

$$P(x, t + \Delta t) = P(x, t) + \frac{\Delta t}{1!} P'(x, t) + \frac{(\Delta t)^2}{2!} P''(x, t) + \frac{(\Delta t)^3}{3!} P'''(x, t) + \dots$$

By solving for the first derivative, we get the following approximation:

$$P'(x, t) = \frac{P(x, t + \Delta t) - P(x, t)}{\Delta t} - \frac{(\Delta t)}{2} P''(x, t) - \dots$$

or, employing the index system:

$$\left(\frac{\partial P}{\partial t}\right)_i^t = \frac{P_i^{t+\Delta t} - P_i^t}{\Delta t} + O(\Delta t).$$

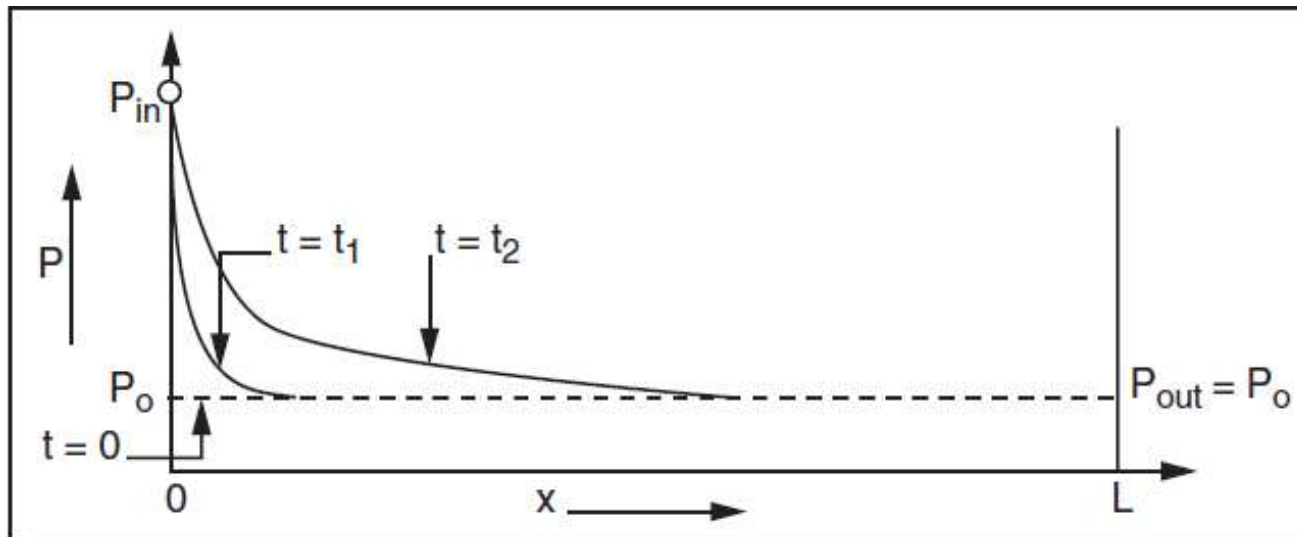
- Here, the error term is proportional to Δt , or of the first order.
- The error therefore approaches zero slower in this case than for the second order term $O(\Delta x^2)$.

Application of Finite Differences to Partial Differential Equations (PDEs)

As an example of PDE, we will take the simplified pressure equation:

$$\frac{\partial^2 p}{\partial x^2} = \frac{1}{\eta} \frac{\partial p}{\partial t} \quad \text{----- (1), where} \quad \eta = \frac{k}{\phi c_t \mu}$$

This is the pressure equation for a 1D system where $0 \leq x \leq L$, where L is the length of the system. After the system is held constant at $P = P_o$, the inlet pressure is raised (at $x = 0$) instantly to $P = P_{in}$ while the outlet pressure is held at $P_{out} = P_o$.



Explicit Finite Difference Approximation of the Linear Pressure Equation

PDE:
$$\frac{\partial^2 p}{\partial x^2} = \frac{1}{\eta} \frac{\partial p}{\partial t} \quad \text{----- (1)}$$

➤ apply finite differences to equation 1 to obtain:

$$\left(\frac{\partial P}{\partial t} \right)_i \approx \frac{P_i^{n+1} - P_i^n}{\Delta t} \quad \text{----- (2)}$$

$$\left(\frac{\partial^2 P}{\partial x^2} \right)_i \approx \frac{P_{i+1}^n + P_{i-1}^n - 2P_i^n}{\Delta x^2} \quad \text{----- (3)}$$

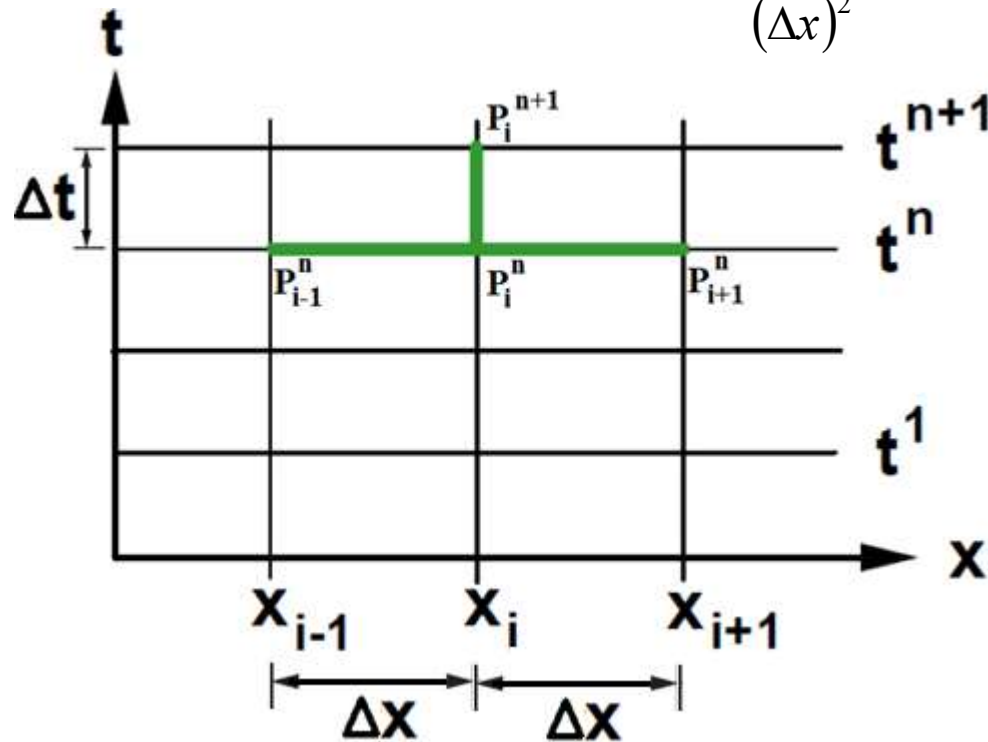
Substituting equations 2 and 3 in equation 1 gives:

$$\text{Explicit Finite Difference } \frac{p_{i+1}^n - 2p_i^n + p_{i-1}^n}{(\Delta x)^2} = \frac{1}{\eta} \left(\frac{p_i^{n+1} - p_i^n}{\Delta t} \right) \text{ ----- (4)}$$

Rearrange equation 4 to obtain an **explicit expression** for, p_i^{n+1} , the only unknown in equation 4:

$$p_i^{n+1} = \alpha p_{i-1}^n + (1 - 2\alpha)p_i^n + \alpha p_{i+1}^n \text{ ----- (5)}$$

$$\text{where } \alpha = \frac{\eta \Delta t}{(\Delta x)^2}$$



Schematic of the explicit finite difference algorithm for solving the simple PDE.

Implicit Finite Difference Approximation of the Linear Pressure Equation

We now return to the original PDE

$$\frac{\partial^2 p}{\partial x^2} = \frac{1}{\eta} \frac{\partial p}{\partial t} \quad \text{----- (1)}$$

The finite difference equation for this case is as follows:

The time derivative is the same as for the explicit finite difference , i.e.

$$\left(\frac{\partial P}{\partial t} \right)_i \approx \frac{P_i^{n+1} - P_i^n}{\Delta t} \quad \text{----- (2)}$$

but the spatial derivative of equation 3 (Part 1) now becomes:

$$\left(\frac{\partial^2 P}{\partial x^2} \right)_i \approx \frac{P_{i+1}^{n+1} + P_{i-1}^{n+1} - 2P_i^{n+1}}{\Delta x^2} \quad \text{----- (6)}$$

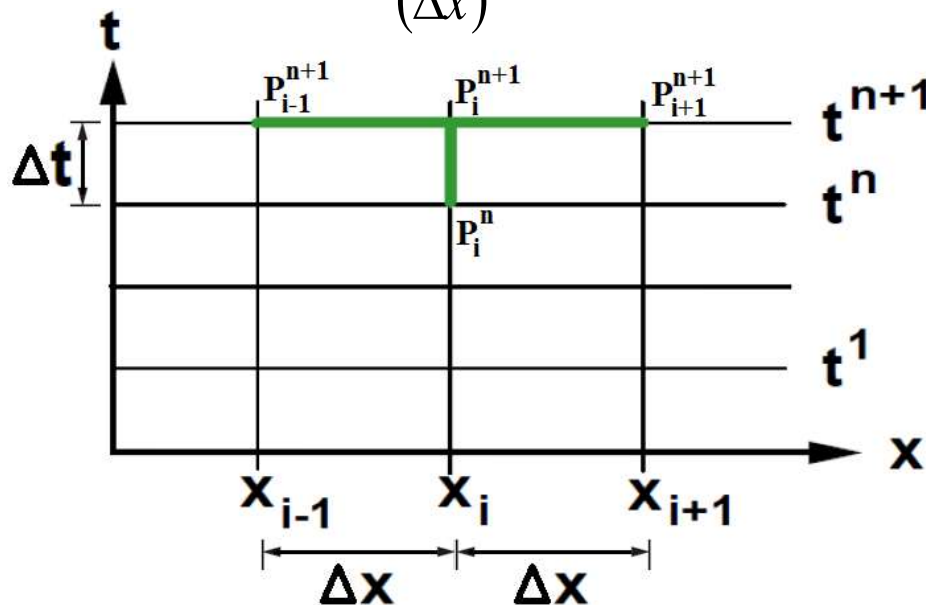
Substituting equations 2 and 6 in equation 1 gives:

$$\text{Implicit Finite Difference } \frac{p_{i+1}^{n+1} - 2p_i^{n+1} + p_{i-1}^{n+1}}{(\Delta x)^2} = \frac{1}{\eta} \left(\frac{p_i^{n+1} - p_i^n}{\Delta t} \right) \text{ ----- (7)}$$

There are three unknowns for each i ; p_{i-1}^{n+1} , p_i^{n+1} , and p_{i+1}^{n+1} . Equation 7 can be rearranged as:

$$-\alpha p_{i-1}^{n+1} + (1 + 2\alpha) p_i^{n+1} - \alpha p_{i+1}^{n+1} = p_i^n \text{ ----- (8)}$$

$$\text{where } \alpha = \frac{\eta \Delta t}{(\Delta x)^2}$$



Schematic of the implicit finite difference algorithm for solving the simple pressure PDE

Discretization

- Converts continuous PDEs into difference form
- Replaces original problem with other problem in terms of algebraic equation which can be solved “easily”
- The reservoir (spatial) domain is represented by spatially distributed, interconnected discrete elements (***grid blocks***)
- Temporal (time) domain is also discretized (***time steps***).

THANK YOU